

Real-time Evolution of $U(1)$ Chiral Charge

Daniel Figueroa
Adrien Florio
Mikhail Shaposhnikov



Lattice 2018, 23th of July 2018

Overview

Lattice Model

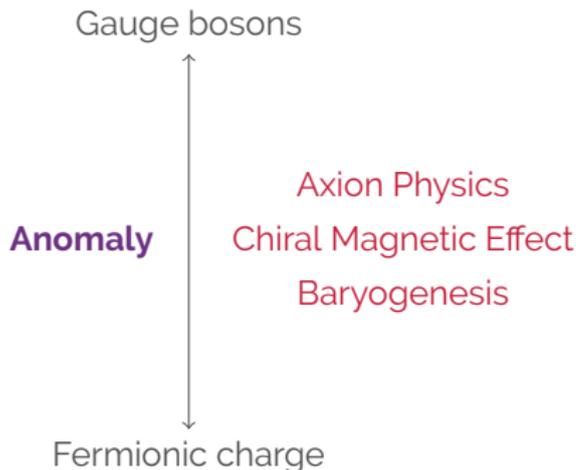
Results/Outlooks

Overview

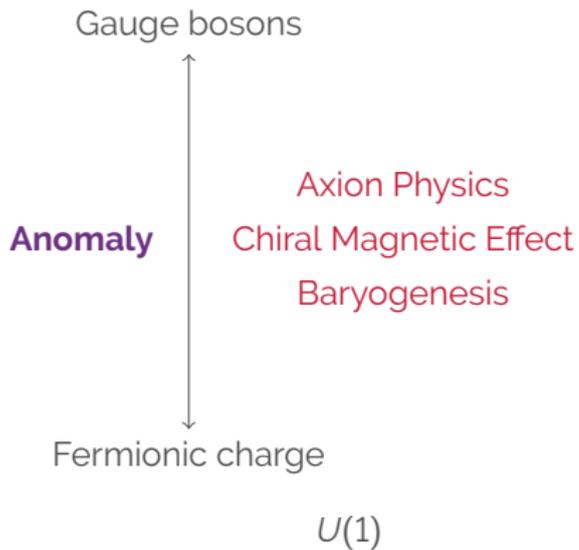
Lattice Model

Results/Outlooks

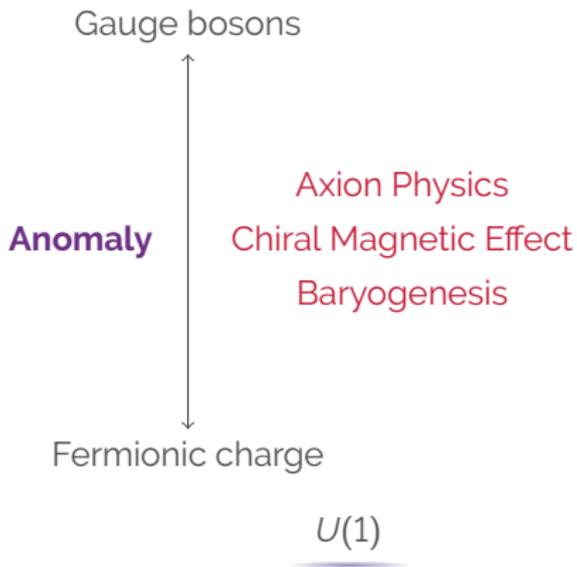
Anomalous Processes



Anomalous Processes



Anomalous Processes



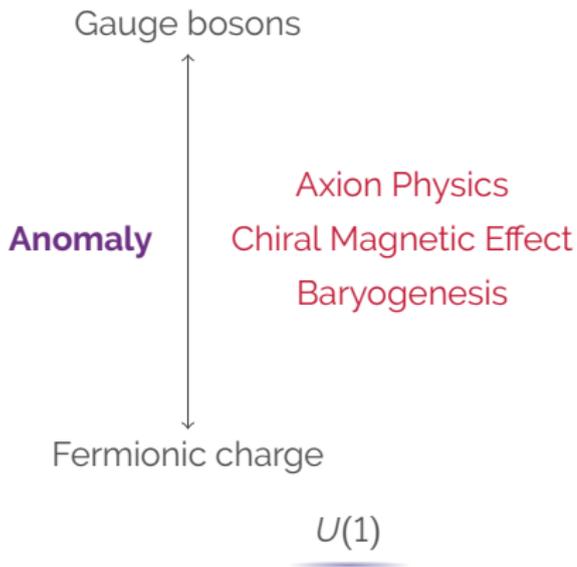
Model

$$\mathcal{L} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} + \bar{\Psi}\not{D}\Psi + (D_{\mu}\phi)^*D_{\mu}\phi - V(\phi)$$

with: $j_5^{\mu} = \bar{\Psi}\gamma^{\mu}\gamma^5\Psi$

$$V(\phi) = m^2|\phi|^2 + \lambda|\phi|^4$$

Anomalous Processes



Model

$$\mathcal{L} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} + \bar{\Psi}\not{D}\Psi + (D_{\mu}\phi)^*D_{\mu}\phi - V(\phi)$$

with: $j_5^{\mu} = \bar{\Psi}\gamma^{\mu}\gamma^5\Psi$

$$V(\phi) = m^2|\phi|^2 + \lambda|\phi|^4$$

Anomaly: $\partial_{\mu}j_5^{\mu} = \frac{e^2}{8\pi^2}F_{\mu\nu}\tilde{F}^{\mu\nu}$

Model

$$\mathcal{L} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} + \bar{\Psi}\not{D}\Psi \\ + (D_{\mu}\phi)^*D_{\mu}\phi - V(\phi)$$

with: $j_5^{\mu} = \bar{\Psi}\gamma^{\mu}\gamma^5\Psi$

$$V(\phi) = m^2|\phi|^2 + \lambda|\phi|^4$$

Anomaly: $\partial_{\mu}j_5^{\mu} = \frac{e^2}{8\pi^2}F_{\mu\nu}\tilde{F}^{\mu\nu}$

Abelian Instabilities

Integrate out $\Psi \longrightarrow F_{CS} = \mu N_{CS}$

Model

$$\mathcal{L} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} + \bar{\Psi}\not{D}\Psi \\ + (D_\mu\phi)^*D_\mu\phi - V(\phi)$$

with: $j_5^\mu = \bar{\Psi}\gamma^\mu\gamma^5\Psi$

$$V(\phi) = m^2|\phi|^2 + \lambda|\phi|^4$$

Anomaly: $\partial_\mu j_5^\mu = \frac{e^2}{8\pi^2}F_{\mu\nu}\tilde{F}^{\mu\nu}$

Abelian Instabilities

$$\frac{e^2}{16\pi^2}F_{\mu\nu}\tilde{F}^{\mu\nu} = \partial_\mu K^\mu$$

$$\int d\vec{x}^3 K^0 = N_{CS} = \frac{\alpha}{2\pi} \int d\vec{x}^3 \vec{A} \cdot \vec{B}$$

Integrate out $\Psi \longrightarrow F_{CS} = \mu N_{CS}$

Model

$$\mathcal{L} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} + \bar{\Psi}\not{D}\Psi \\ + (D_{\mu}\phi)^*D_{\mu}\phi - V(\phi)$$

with: $j_5^{\mu} = \bar{\Psi}\gamma^{\mu}\gamma^5\Psi$
 $V(\phi) = m^2|\phi|^2 + \lambda|\phi|^4$

Anomaly: $\partial_{\mu}j_5^{\mu} = \frac{e^2}{8\pi^2}F_{\mu\nu}\tilde{F}^{\mu\nu}$

Abelian Instabilities

$$\frac{e^2}{16\pi^2}F_{\mu\nu}\tilde{F}^{\mu\nu} = \partial_{\mu}K^{\mu}$$

$$\int d\vec{x}^3 K^0 = N_{CS} = \frac{\alpha}{2\pi} \int d\vec{x}^3 \vec{A} \cdot \vec{B}$$

Integrate out $\Psi \longrightarrow F_{CS} = \mu N_{CS}$

$k^2 AA$ vs $\mu k AA$

Model

$$\mathcal{L} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} + \bar{\Psi}\not{D}\Psi \\ + (D_\mu\phi)^*D_\mu\phi - V(\phi)$$

with: $j_5^\mu = \bar{\Psi}\gamma^\mu\gamma^5\Psi$
 $V(\phi) = m^2|\phi|^2 + \lambda|\phi|^4$

Anomaly: $\partial_\mu j_5^\mu = \frac{e^2}{8\pi^2}F_{\mu\nu}\tilde{F}^{\mu\nu}$

Abelian Instabilities

$$\frac{e^2}{16\pi^2}F_{\mu\nu}\tilde{F}^{\mu\nu} = \partial_\mu K^\mu$$

$$\int d\vec{x}^3 K^0 = N_{CS} = \frac{\alpha}{2\pi} \int d\vec{x}^3 \vec{A} \cdot \vec{B}$$

Integrate out $\Psi \longrightarrow F_{CS} = \mu N_{CS}$

$k^2 AA$ vs $\mu k AA$

Instability if $k < \frac{\alpha}{\pi}\mu!$

Abelian Instabilities

$$\frac{e^2}{16\pi^2} F_{\mu\nu} \tilde{F}^{\mu\nu} = \partial_\mu K^\mu$$

$$\int d\vec{x}^3 K^0 = N_{CS} = \frac{\alpha}{2\pi} \int d\vec{x}^3 \vec{A} \cdot \vec{B}$$

Integrate out $\Psi \longrightarrow F_{CS} = \mu N_{CS}$

$k^2 AA$ vs $\mu k AA$

Instability if $k < \frac{\alpha}{\pi} \mu!$

Comments

- Long-range gauge fields
- Symmetric phase
- Also at non-zero temperature

Comments

- Long-range gauge fields
- Symmetric phase
- Also at non-zero temperature

External Magnetic Field

EoM + Flux Cons. $\implies \vec{B}$ vac. state

Comments

- Long-range gauge fields
- Symmetric phase
- Also at non-zero temperature

External Magnetic Field

EoM + Flux Cons. $\implies \vec{B}$ vac. state

Generating \vec{A} cost no energy

Comments

- Long-range gauge fields
- Symmetric phase
- Also at non-zero temperature

External Magnetic Field

EoM + Flux Cons. $\implies \vec{B}$ vac. state



Continuum of vac. $N_{CS} \propto \vec{A} \cdot \vec{B}$



Generating \vec{A} cost no energy

External Magnetic Field

EoM + Flux Cons. $\implies \vec{B}$ vac. state



Continuum of vac. $N_{CS} \propto \vec{A} \cdot \vec{B}$



Generating \vec{A} cost no energy

Sum-Up

- $B = 0$: Instability
- $B \neq 0$: Non-trivial vac. structure



Similar to non-abelian!

External Magnetic Field

EoM + Flux Cons. $\implies \vec{B}$ vac. state



Continuum of vac. $N_{CS} \propto \vec{A} \cdot \vec{B}$



Generating \vec{A} cost no energy

Sum-Up

- $B = 0$: Instability
- $B \neq 0$: Non-trivial vac. structure



Similar to non-abelian!

Overview

Lattice Model

Results/Outlooks

Lattice Model

$$\mathcal{L} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} + (D_\mu\phi)^*D_\mu\phi - V(\phi) \\ -\frac{1}{2c_s^2}(\partial_0 a)^2 + \frac{1}{2}(\partial_i a)^2$$

with a a scalar field

*U(1)*Axion!*

Reproduce EoM

Real-Time Simulations

MC generated thermal ensemble

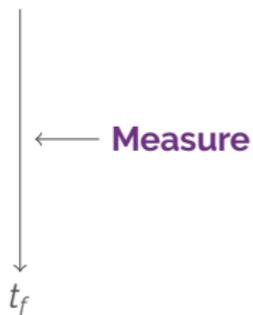
← **Measure**

t_f

Solve **classical** EoM

Real-Time Simulations

MC generated thermal ensemble



Solve **classical** EoM

Technical Comments

- Lattice topological $F\tilde{F}$
- External mag. field as twisted BCs
- Homogeneous axion much easier

Refs.: [JHEP04\(2018\)026](#)
[j.nuclphysb.2017.12.001](#)

Overview

Lattice Model

Results/Outlooks

Chern-Simons Density

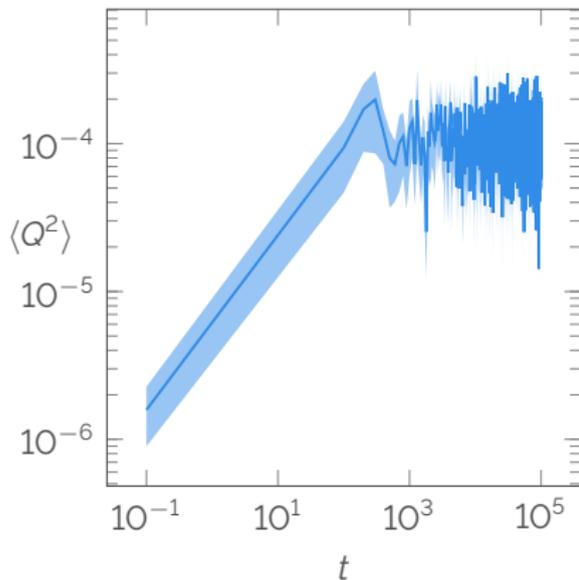
- $\mu = 0, B = 0$

$$\vec{A} \cdot \vec{B} \text{ costs } E \implies \langle Q^2(t) \rangle \rightarrow cst$$

- $\mu = 0, B \neq 0$

$$\vec{A} \cdot \vec{B} \text{ costs } E \implies \langle Q^2(t) \rangle \xrightarrow{\text{RD walk}} \Gamma Vt$$

CS Evolution



Chern-Simons Density

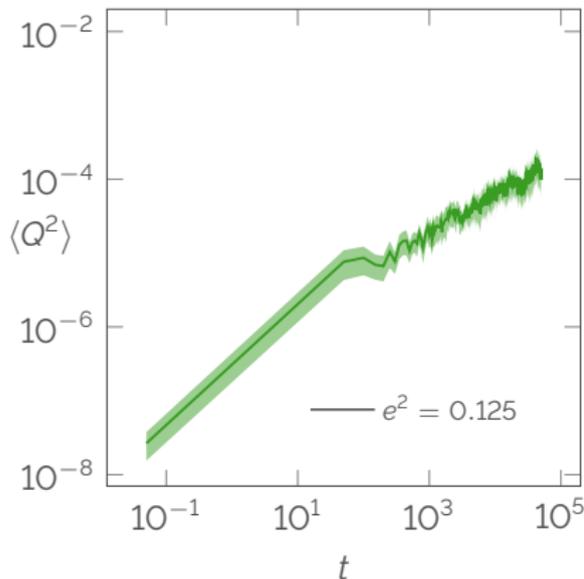
- $\mu = 0, B = 0$

$$\vec{A} \cdot \vec{B} \text{ costs } E \implies \langle Q^2(t) \rangle \rightarrow cst$$

- $\mu = 0, B \neq 0$

$$\vec{A} \cdot \vec{B} \text{ costs } E \implies \langle Q^2(t) \rangle \xrightarrow{\text{RD walk}} \Gamma Vt$$

CS Evolution



Chern-Simons Density

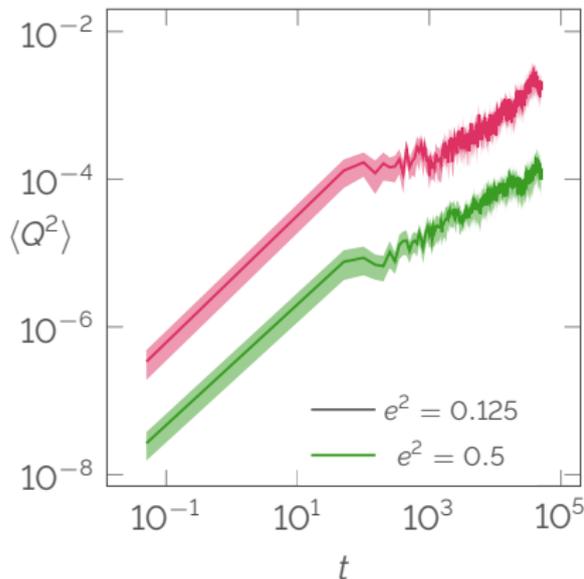
- $\mu = 0, B = 0$

$$\vec{A} \cdot \vec{B} \text{ costs } E \implies \langle Q^2(t) \rangle \rightarrow cst$$

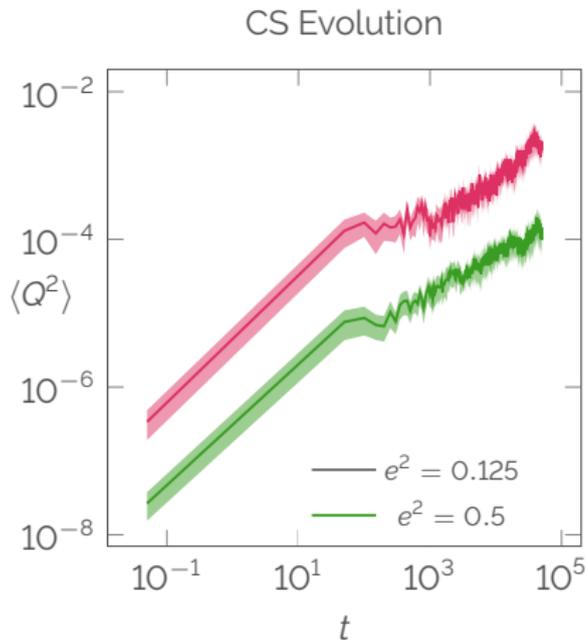
- $\mu = 0, B \neq 0$

$$\vec{A} \cdot \vec{B} \text{ costs } E \implies \langle Q^2(t) \rangle \xrightarrow{\text{RD walk}} \Gamma \sqrt{t}$$

CS Evolution



Results



$$\Gamma \text{ pred. by MHD: } \frac{\Gamma_{th}}{e^6 B^2} \approx 2.5 \cdot 10^{-5}$$

$$\text{Measured } \Gamma: \frac{\Gamma_{exp}}{e^6 B^2} \approx 1.5 \pm 0.2 \cdot 10^{-3}$$

Agree on parametric

but

$$\frac{\Gamma_{exp}}{\Gamma_{th}} \approx 60!$$

Results

Γ pred. by MHD: $\frac{\Gamma_{th}}{e^6 B^2} \approx 2.5 \cdot 10^{-5}$

Measured Γ : $\frac{\Gamma_{exp}}{e^6 B^2} \approx 1.5 \pm 0.2 \cdot 10^{-3}$

Agree on parametric

but

$$\frac{\Gamma_{exp}}{\Gamma_{th}} \approx 60!$$

Next: Chemical Potential

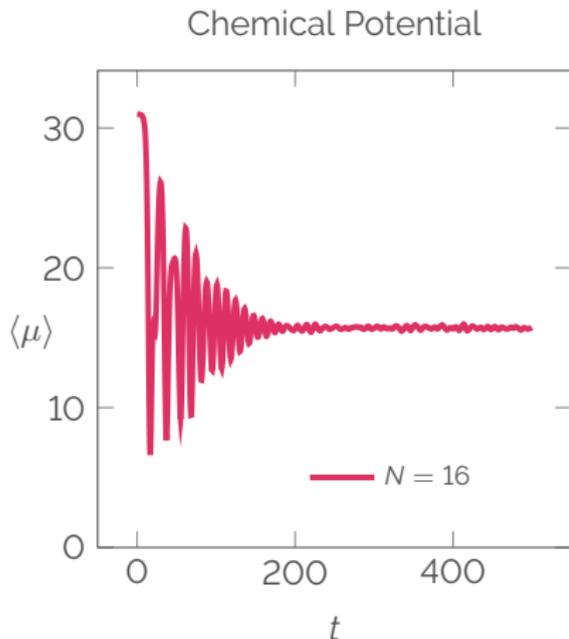
- $\mu \neq 0, B = 0$

Lat. art.: $\mu_{min} = k_{min} \frac{\pi}{\alpha} \propto \frac{1}{N}$

Next: Chemical Potential

- $\mu \neq 0, B = 0$

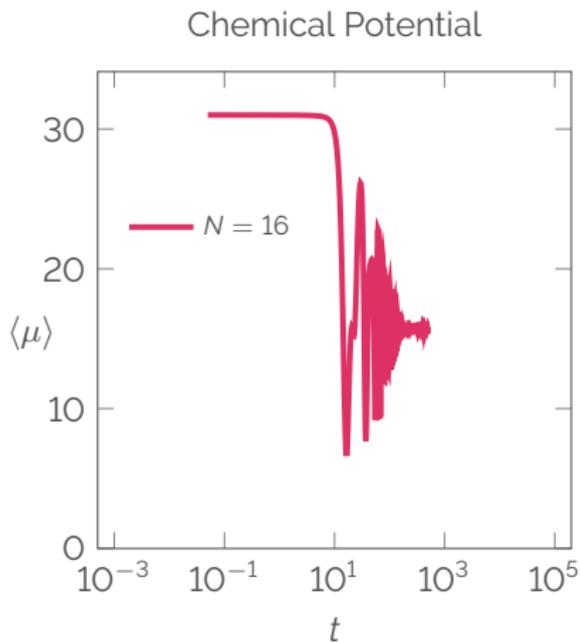
Lat. art.: $\mu_{min} = k_{min} \frac{\pi}{\alpha} \propto \frac{1}{N}$



Next: Chemical Potential

- $\mu \neq 0, B = 0$

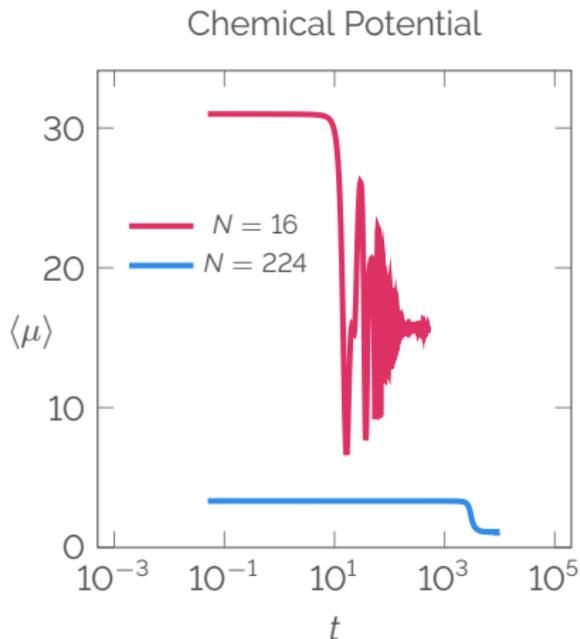
Lat. art.: $\mu_{min} = k_{min} \frac{\pi}{\alpha} \propto \frac{1}{N}$



Next: Chemical Potential

- $\mu \neq 0, B = 0$

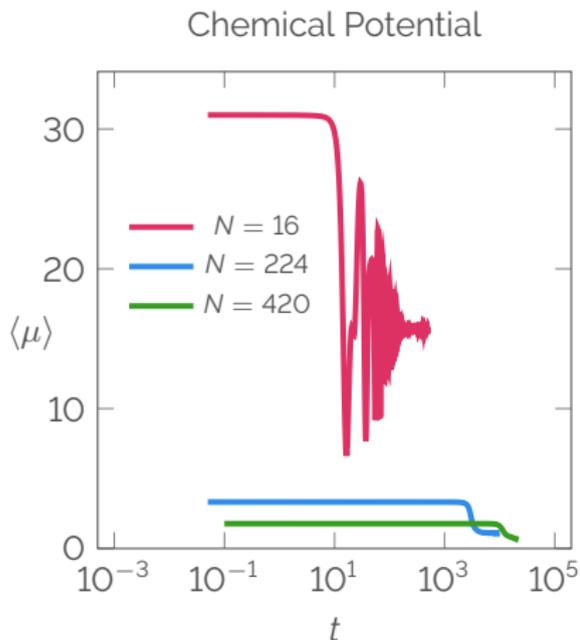
Lat. art.: $\mu_{min} = k_{min} \frac{\pi}{\alpha} \propto \frac{1}{N}$



Next: Chemical Potential

- $\mu \neq 0, B = 0$

Lat. art.: $\mu_{min} = k_{min} \frac{\pi}{\alpha} \propto \frac{1}{N}$



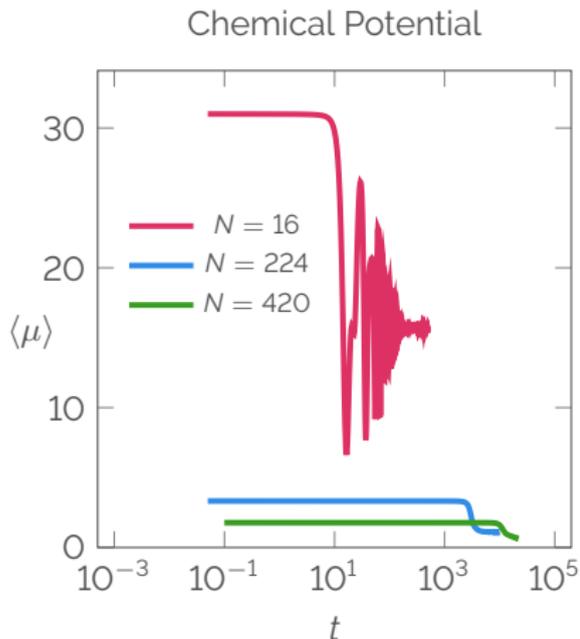
Next: Chemical Potential

- $\mu \neq 0, B = 0$

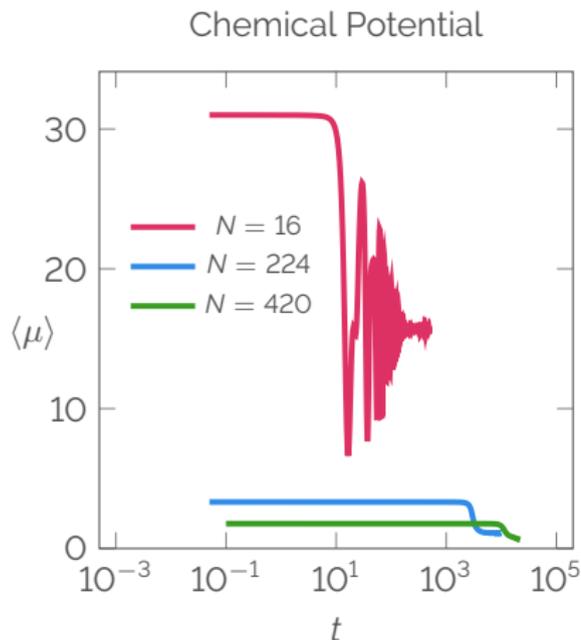
Lat. art.: $\mu_{min} = k_{min} \frac{\pi}{\alpha} \propto \frac{1}{N}$

Question:

μ independent rate at small μ ?

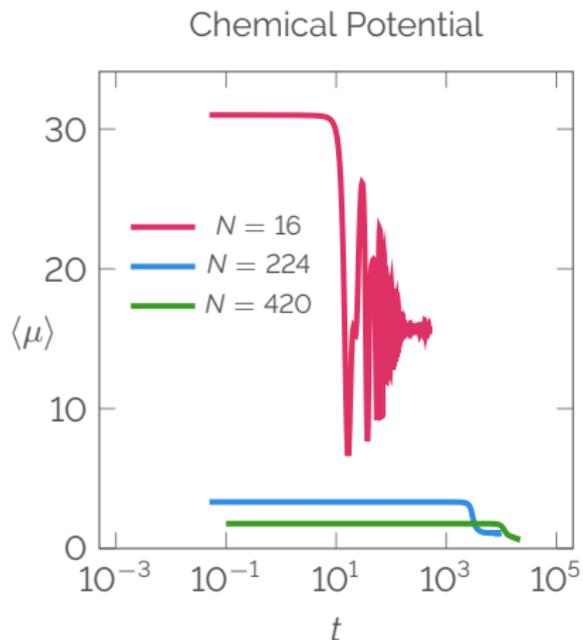


Further Outlooks



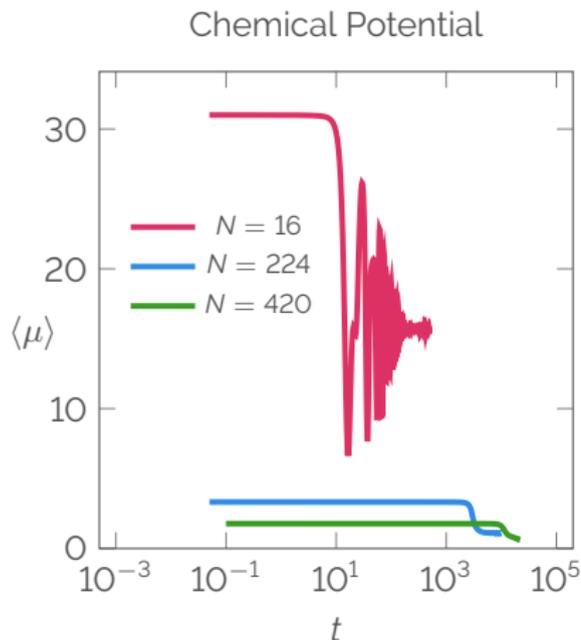
- $B \neq 0, \mu \neq 0$
- Hard Thermal Loops
- Non-homogeneous axion

Further Outlooks



- $B \neq 0, \mu \neq 0$
- Hard Thermal Loops
- Non-homogeneous axion

Further Outlooks



- $B \neq 0, \mu \neq 0$
- Hard Thermal Loops
- Non-homogeneous axion

Further Outlooks

- $B \neq 0, \mu \neq 0$
- Hard Thermal Loops
- Non-homogeneous axion

Take Away

- **Discrepancies between MHD and full simulations**
- Exciting outlooks

Further Outlooks

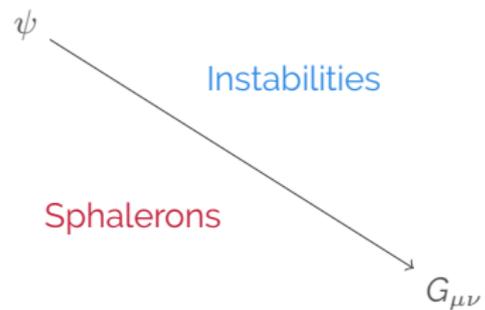
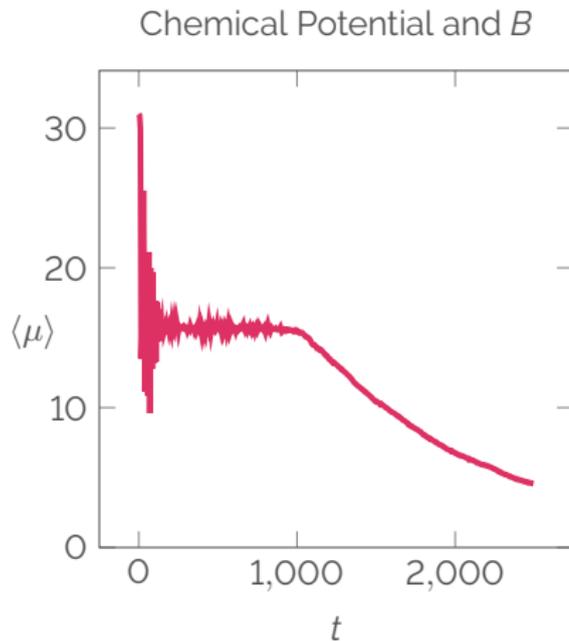
- $B \neq 0, \mu \neq 0$
- Hard Thermal Loops
- Non-homogeneous axion

Take Away

- Discrepancies between MHD and full simulations
- **Exciting outlooks**

Thank you!

Non-Abelian



[Rubakov,1986]

[Rummukainen,2014]